

## PROPAGATIONS FROM A SPACE CURVE IN THREE SPACE WITH INDICATRIX A SURFACE

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ABSTRACT. Generic singularities of rays emanating from a space curve in  $\mathbb{R}^3$  in all directions with the rate determined by an indicatrix (independent of the point in  $\mathbb{R}^3$ ) defined by a surface are classified. Similarly rays emanating from surface defined by an indicatrix given by a curve are also considered. Some applications to control theory are indicated.

### 1. INTRODUCTION

In this paper we solve two problems on the classification of local geometrical singularities that are related to control theory. We use some techniques from the singularity theory of caustics and wave fronts to study singularities of exponential mappings in a class of control problems which correspond to special integrable Hamiltonian systems with straight lines as extremals.

The first problem concerns a control system on a three-dimensional affine space with points  $q \in \mathbb{R}^3$ . We identify the tangent space  $T_q\mathbb{R}^3$  with  $\mathbb{R}^3$  itself. At each point  $q$  we choose an indicatrix  $I_q$  of admissible velocities  $\dot{q} = \frac{\partial q}{\partial \mu}$  of motion which we assume is independent of the point  $q$  itself. Assume that this set is parametrised locally by a regular mapping  $(x, y) \mapsto r_2(x, y)$  whose image is a surface  $M$ . We shall now write  $M$  in place of  $I_q$ .

An *admissible motion* is a smooth curve  $\gamma(\mu) \in \mathbb{R}^3$ , parametrised by a segment of the (affine) time axis  $\mu$ , such that the velocity at each point  $\dot{\gamma}$  belongs to the set of admissible velocities  $M$ .

Let  $q_b(\mu)$  be the trajectory of an admissible motion of an initial point  $b \in N$ , issuing at  $\mu = 0$  from an initial set  $N$ , where  $N$  is a space curve which is a submanifold in  $\mathbb{R}^3$ .

For a fixed value  $\mu = \mu_0$  let  $\mathcal{C}_b$  be the Banach manifold of all admissible trajectories defined on the segment  $[0, \mu_0]$  from an initial point  $b \in N$ .

Consider the endpoint mapping  $\mathcal{E}_b : \mathcal{C}_b \rightarrow \mathbb{R}^3$  which associates the endpoint  $q_b(\mu_0)$  to a trajectory  $q_b(\mu)$ .

A corollary of the Pontryagin maximum principle, see [1, 9], states that critical values of  $\mathcal{E}$  for all  $\mu_0$  trace extremal trajectories. In our case these are projections to  $\mathbb{R}^3$  of solutions of the associated Hamiltonian canonical equations on the cotangent bundle

$$\dot{q} = \frac{\partial H_*(p, q)}{\partial p}, \quad \dot{p} = -\frac{\partial H_*(p, q)}{\partial q}.$$

Here the Hamiltonian function  $H_*(p, q)$  on the cotangent bundle  $T^*\mathbb{R}^3$  is the restriction (multivalued in general) to the subset  $\{(p, q) \mid \exists(x, y) : \frac{\partial H(p, q, x, y)}{\partial x} = \frac{\partial H(p, q, x, y)}{\partial y} = 0\}$  of the function  $H = \langle p, r_2(x, y) \rangle$ , provided that the initial conditions  $(p_0, q_0)$  satisfy the relation  $\langle p_0, v \rangle = 0$  for each vector  $v$  tangent to  $N$  at  $b$ . The angle brackets  $\langle -, - \rangle$  denote the standard pairing of vectors  $\mathbb{R}^3$  and covectors  $p$  of the dual space  $(\mathbb{R}^3)^\wedge$ .

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In our case extremals are straight lines (parametrised by  $\mu \in \mathbb{R}$ )  $b + \mu v$ ,  $b \in N$ ,  $v \in M$ , such that there is a covector  $p_0$ , which annihilates both tangent spaces  $T_b N$  and  $T_v M$ . Points on these lines with fixed  $\mu$  form a wave front  $E_\mu$  of the Legendre variety  $L_\mu$ , which is the image of the Legendre submanifold  $L_0$  of the initial conditions under the Hamiltonian flow.

The envelope  $B(N, M)$  of these extremals is the union of singular points of sets of critical values of  $\mathcal{E}_b$  for all  $b$  and  $\mu_0$ .

Rephrasing the above in physical terms, consider an initial space curve in three space which emits rays from every point, and such that the speed of a ray at any point is completely determined by its direction. The *boundary of the set of attainability* of the rays after a given time  $\mu$  will be the wave front  $E_\mu$ . The caustics or focal sets correspond to the singularities of this set of attainability.

In this first problem as described above we consider an initial space curve with a velocity indicatrix defined by a surface. The classification where the indicatrix was also a space curve (independent of the point in  $\mathbb{R}^3$ ) was given in [8]. The case of an initial surface and a velocity indicatrix described by a surface was studied in [3]. For completeness we also consider in the present paper a second problem interchanging the surface and the curve, i.e. the indicatrix is defined by a space curve and the initial manifold is a surface.

At present this second problem seems to have fewer applications than the first despite the fact that away from the initial surface the classification coincides with that of the first problem.

In the first problem the dimension of the indicatrix  $M$  is one less than the dimension of the ambient space  $\mathbb{R}^3$ , so the wave fronts  $E_\mu$  form a family of equidistants in Finsler geometry. However as the dimension of the indicatrix in the second problem is not one less than the ambient space the wave fronts do not form a family of equidistants in Finsler geometry.

In this paper we classify the possible generic singularities of the envelopes  $B(M, N)$  and of the family of wave fronts in both problems. We also classify the generic singularities near the initial surface itself in the second problem.

The method of classification of the singularities is similar to that of a related problem in [7]. In that paper the wave fronts were taken to be the closure of an affine ratio of pairs of points, one from a curve and the other from a surface that share parallel tangent planes. Here we consider in the first problem the surface, and then in the second problem the curve, to be at infinity. The computations were largely omitted from [7] and since in the present context they are slightly easier to write down we take this opportunity to include more details.

**1.1. Main definitions and results.** Let  $M$  be a smooth surface and let  $N$  be a smooth space curve both embedded in affine three space.

Consider a pair  $a, b$  of points  $a \in M$  and  $b \in N$  such that the plane tangent to the surface  $M$  at  $a$  is parallel to some plane tangent to  $N$  at  $b$ . The pair  $a, b$  is called a *parallel pair* and the straight line through  $a, b$  is called a *chord*. The envelope of the family of all chords is called the Minkowski set of  $M$  and  $N$ . In this paper we shall classify its generic singularities.

The chord  $l(a, b)$  joining the parallel pair is defined by

$$(1) \quad l(a, b) = \{q \in \mathbb{R}^3 \mid q = \mu a + (1 - \mu)b, \mu \in \mathbb{R}\}.$$

The following definitions are valid for the propagating from the space curve case but similar definitions, by replacing  $\mu$  with  $1 - \mu = \lambda$  hold in the propagating from the surface case. The points which correspond to parallel tangent plane is the furthest point of the wave front from the curve. The wave front  $E_\mu$  is the boundary of where the rays have reached at time  $\mu$ . In the previous papers [3, 4, 7] barycentric coordinates were introduced on to the chords. Here however, we omit  $\lambda$  and just use the coordinate  $\mu$  where  $\mu = 0$  corresponds to the point on the curve  $N$  and  $\mu = \infty$  corresponds to the point on the surface  $M$ .

A germ of the *affine  $\mu$ -equidistant*  $E_\mu$  of the pair  $(M, N)$  is the set of points  $q \in \mathbb{R}^3$  such that  $q = \mu a + b$  for given  $\mu \in \mathbb{R}$  and for all parallel pairs  $(a, b)$  close to  $(a_0, b_0)$ . Note that  $E_0$  is the germ of  $N$  at  $b_0$ .

The space  $\mathbb{R}_e^4 = \mathbb{R} \times \mathbb{R}^3$  with coordinate  $\mu \in \mathbb{R}$  (affine time), on the first factor is called the *extended affine space*. Denote by

$\rho : (\mu, q) \mapsto \mu$  the projection of  $\mathbb{R}_e^4$  to the first factor and by  $\pi : (\mu, q) \mapsto q$  the projection to the second factor.

The *affine extended wave front*  $W(M, N)$  of the pair  $M, N$  is the union of all affine equidistants each embedded into its own slice of the extended affine space:

$$W(M, N) = \{(\mu, E_\mu)\} \subset \mathbb{R}_e^4.$$

The *bifurcation set*  $B(M, N)$  of a family of affine equidistants (or of the family of chords) of the pair  $(M, N)$  is the image under  $\pi$  of the locus of the critical points of the restriction  $\pi_r = \pi|_{W(M, N)}$ . A point is *critical* if  $\pi_r$  at this point fails to be a regular projection of a smooth submanifold. In general  $B(M, N)$  consists of two components: the *caustic*  $\Sigma$  is the projection of the singular locus of the extended wave front  $W(M, N)$  and the *criminant*  $\Delta$  is the (closure of the) image under  $\pi_r$  of the set of regular points of  $W(M, N)$  which are the critical points of the projection  $\pi$  restricted to the regular part of  $W(M, N)$ . The caustic consists of the singular points of the momentary equidistants  $E_\mu$  while the criminant is the envelope of the family of regular parts of the momentary equidistants. Besides being swept out by the momentary equidistants, the affine extended wave front is swept out by the liftings to  $\mathbb{R}_e^{3+1}$  of chords. Each of them has a regular projection to the configuration space  $\mathbb{R}^3$ . Hence the bifurcation set  $B(M, N)$  is essentially the envelope of the family of chords.

In the generic setting the distinguished chords split into three distinct sub-cases: In the first (*transversal*) case the base points  $a_0 \in M$  and  $b_0 \in N$  are distinct and the chord through them is transversal to both  $M$  and  $N$ . In the second (*tangential*) case the base points  $a_0 \in M$  and  $b_0 \in N$  are distinct but the tangent line to  $N$  lies in the tangent plane to  $M$ . A subcase of the tangential case called the *supertangential* case occurs when the line tangent to the curve  $N$  at  $b$  contains the point  $a$ , i.e. the chord and the tangent line are the same.

**Definition 1.1.** Two germs of families  $F_1$  and  $F_2$  in parameters  $\mu, q$  are called *space-time contact equivalent* if there exists a nonzero function  $\phi(z, \mu, q)$  and diffeomorphism  $\hat{\theta} : (z, \mu, q) \mapsto (Z(z, \mu, q), P(\mu, q), Q(q))$  such that  $F_1 = \phi F_2 \circ \hat{\theta}$ .

Notice that the diffeomorphism  $\hat{\theta} : (\mu, q) \mapsto (P(\mu, q), Q(q))$  of the total parameter space  $\mathbb{R}^{3+1}$  maps the extended wave front of the first family to the extended wave front of the second family and the diffeomorphism  $\hat{\theta} : q \mapsto Q(q)$  of the  $q$ -parameter space  $\mathbb{R}^3$  maps the bifurcation set of the first family to the bifurcation set of the second family.

**Definition 1.2.** Two germs of families  $F_1$  and  $F_2$  are called *time-space contact equivalent* if there exists a nonzero function  $\phi(z, \mu, q)$  and diffeomorphism  $\tilde{\theta} : (z, \mu, q) \mapsto (Z(z, \mu, q), P(\mu), Q(\mu, q))$  such that  $F_1 = \phi F_2 \circ \tilde{\theta}$ .

The diffeomorphism  $\tilde{\theta} : (\mu, q) \mapsto (P(\mu), Q(\mu, q))$  preserves the fibration of the  $\mu, q$  space into fibres parallel to the  $q$  space. If two families are time-space contact equivalent then their respective families of momentary wave fronts are diffeomorphic.

The main results are as follows: The first theorem which concerns the wave fronts follows immediately from the results of [7].

**Theorem 1.3.** *The families of wave fronts and their bifurcations in the propagating from the curve and in the propagating from the surface cases are diffeomorphic to those in the affine ratio case [7]. In fact the generating functions are time-space contact equivalent.*

The following theorems all concern the projection  $\pi$  and are related to the caustics. The theorems are the complete classification of generic singularities in the various settings. The list is the same as in [7]. Unlike theorem 1.3 they do not follow immediately from the previous papers and require separate calculation.

**Theorem 1.4.** *In the propagating from the curve transversal case outside  $M$  and  $N$  the germ at any point of the envelope of the family of chords for generic  $M$  and  $N$  is diffeomorphic to one of the standard caustics of  $A_k$  type with  $k = 2, 3$  or  $4$  (regular surface, cuspidal edge or swallowtail).*

**Theorem 1.5.** *In the propagating from the curve tangential case the germ at any point outside  $M$  and  $N$  of the envelope of the family of chords for generic  $N$  and  $M$  is diffeomorphic to one of the standard caustics of the boundary singularities of the types  $B_2, B_3, B_4, C_3, C_4$  or  $F_4$ . If moreover the tangent line to the curve coincides with the chord (supertangential case) then generically only  $B_2$  and  $C_3$  occur.*

**Theorem 1.6.** *In the propagating from the surface transversal case outside  $M$  and  $N$  the germ at any point of the envelope of the family of chords for generic  $M$  and  $N$  is diffeomorphic to one of the standard caustics of  $A_k$  type with  $k = 2, 3$  or  $4$  (regular surface, cuspidal edge or swallowtail).*

**Theorem 1.7.** *In the propagating from the surface transversal case the envelope of the family of chords transversally intersects the surface  $M$  when  $\lambda = 0$  generically at either its smooth points or at points of a cuspidal ridge.*

**Theorem 1.8.** *In the propagating from the surface tangential setting the germ at any point outside  $M$  and  $N$  of the envelope of the family of chords for generic  $N$  and  $M$  is diffeomorphic to one of the standard caustics of the boundary singularities of the types  $B_2, B_3, B_4, C_3, C_4$  or  $F_4$ . If moreover the tangent line to the curve coincides with the chord (supertangential case) then generically only  $B_2$  and  $C_3$  occur.*

**1.2. Generating families.** Now consider the following generating family  $\mathcal{F}_1$  in the propagating from the curve case. The family has variables  $n \in (\mathbb{R}^3)^\wedge \setminus \{0\}$ ,  $t$  and  $(x, y)$ , and parameters  $(\mu, q) \in \mathbb{R} \times \mathbb{R}^3$ ;

$$(2) \quad \mathcal{F}_1(n, t, x, y, \mu, q) = \langle r_1(t) + \mu r_2(x, y) - q, n \rangle$$

where  $r_1(t)$  is the embedding with the image  $N$ , and  $r_2(x, y)$  is the embedding with image  $M$ .

In the propagating from the surface case we use the generating family

$$(3) \quad \mathcal{F}_2(n, t, x, y, \lambda, q) = \langle \lambda r_1(t) + r_2(x, y) - q, n \rangle$$

with the same variables as  $\mathcal{F}_1$  but parameters  $(\lambda, q) \in \mathbb{R} \times \mathbb{R}^3$ .

In the paper [7] the affine ratio case was studied and the generic bifurcations of the wave fronts were classified. There the generating family used was

$$(4) \quad \tilde{\mathcal{F}}(n, t, x, y, \mu, q) = \langle (1 - \mu)r_1(t) + \mu r_2(x, y) - q, n \rangle.$$

We now show that the two family germs  $\widehat{\mathcal{F}}$  and  $\mathcal{F}_1$  are time-space equivalent (see theorem 1.3) and hence the classification of their generic wave fronts and their bifurcations are in fact the same.

**Proof of theorem 1.3.**

Assuming that  $\mu \neq 1$  we can divide the family (4) by  $(1 - \mu)$  to give

$$\widehat{\mathcal{F}}(n, t, x, y, \mu, q) = \langle r_1(t) + \frac{\mu}{1 - \mu} r_2(x, y) - \frac{q}{1 - \mu}, n \rangle.$$

This is time-space contact equivalent to

$$\widehat{\mathcal{F}}_1 = \langle r_1(t) + \tilde{\mu} r_2(x, y) - \tilde{q}, n \rangle$$

with  $\tilde{\mu} = \frac{\mu}{1 - \mu}$  and  $\tilde{q} = -\frac{q}{1 - \mu}$ .  $\square$

If  $\mu = 1$  then this is equivalent to the present case at infinity and so does not appear in this "non-projective" setting. Similar considerations show that the family  $\widehat{\mathcal{F}}$  is time-space contact equivalent to the family  $\mathcal{F}_2$ .

## 2. PROPAGATING FROM THE CURVE IN THE TRANSVERSAL SETTING

In the transversal setting up to an appropriate affine transformation of  $\mathbb{R}^3$  we can always assume that in some coordinate system  $(x, y, z)$  the base parallel pair  $a_0, b_0$  coincides with the pair of points  $(0, 0, -1), (0, 0, 0)$ , the tangent plane to the surface  $M$  at  $a_0$  is parallel to the  $(x, y)$ -coordinate plane and the tangent line to the curve  $N$  at  $b_0$  coincides with the  $x$ -axis.

In these coordinates the surface  $M$  in the neighbourhood of  $a_0$  is the graph  $M = \{(x, y, z) | z = f(x, y) - 1\}$  of the function  $f$  with vanishing 1-jet. Let  $f(x, y) = \sum_{i+j \geq 2} f_{ij} x^i y^j$  be the Taylor decomposition of the germ of  $f$  at the origin. Define the curve  $N$  to be the set  $\{(t, \alpha(t), \beta(t))\}$  with the functions  $\alpha(t) = \alpha_2 t^2 + \alpha_3 t^3 + \dots$  and  $\beta(t) = \beta_2 t^2 + \beta_3 t^3 + \dots$  starting with at least quadratic terms in  $t$ .

**Proposition 2.1.** *The germ of the family  $\mathcal{F}_1$  at a point corresponding to a point on the base chord is stably-equivalent to the product of the family germ  $\Phi(t, \mu, q) = \beta(t) + \mu[f(\widehat{x}, \widehat{y}) - 1] - q_3$  at the subset  $\widehat{S}_0 = \{t = 0, q_1 = q_2 = 0\}$  with a nonzero factor. Here we use the substitution  $\widehat{x} = \frac{q_1 - t}{\mu}, \widehat{y} = \frac{q_2 - \alpha(t)}{\mu}$ .*

**Proof.** Writing the family  $\mathcal{F}_1$  in the coordinate form we get

$$\mathcal{F}_1 = An_1 + Bn_2 + Cn_3$$

where

$$A = t + \mu x - q_1$$

$$B = \alpha(t) + \mu y - q_2$$

and

$$C = \beta(t) + \mu[f(x, y) - 1] - q_3$$

For  $\mu \neq 0$  the functions  $A$  and  $B$  are regular and we can choose  $A, B$  as the coordinate functions instead of  $x$  and  $y$ . In particular we can write  $x = \frac{A + q_1 - t}{\mu}$  and  $y = \frac{B + q_2 - \alpha(t)}{\mu}$ .

So in the new coordinates we have  $\mathcal{F}_1 = An_1 + Bn_2 + C(A, B, t, \mu, q)n_3$ . The function  $C$  does not depend on  $n_1$  and  $n_2$  and the Hadamard lemma implies  $C(A, B, t, \mu, q) = C(0, 0, \mu, t, q) + A\varphi_1 + B\varphi_2$ , where  $\varphi_1$  and  $\varphi_2$  are smooth functions in  $A, B, t, \mu$  and  $q$  which vanish at the origin.

Now the function  $\mathcal{F}_1$  takes the form  $\mathcal{F}_1 = A(n_1 + \varphi_1 n_3) + B(n_2 + \varphi_2 n_3) + C(0, 0, t, \mu, q)$  where the first two terms represent a non degenerate quadratic form in the independent variables  $A, (n_1 + \varphi_1 n_3), B$  and  $(n_2 + \varphi_2 n_3)$  in the vicinity of the point on the base chord.

Therefore, the function  $\mathcal{F}_1$  is stably-equivalent to the function  $\Phi = C(0, 0, t, \mu, q)$  being the restriction of the function  $C$  to the subspace  $A = B = 0$ . So to study the envelope of chords and the families of wave fronts we can study the family germ

$$\Phi(t, \mu, q) = \beta(t) + \mu \left[ f \left( \frac{q_1 - t}{\mu}, \frac{q_2 - \alpha(t)}{\mu} \right) - 1 \right] - q_3 \quad \square.$$

For the family germ  $\Phi$  at the point  $m_0 = (0, \mu_0, 0, 0, q_3 = -\mu_0)$ , on the base chord  $l(a_0, b_0)$  denote by  $g(t) = \Phi(t, \mu_0, 0, 0, q_3)$  at  $t$  the respective organising centre function.

To determine the singularity type of the generating family germ  $\Phi$  at the point  $m_0$  and the respective versality conditions denote by  $\Phi_k(\mu, q)$  the coefficients at  $t^k$  in the Taylor decomposition of  $\Phi$  with respect to  $t$  at the origin.

$$\Phi = \Phi_0 + \Phi_1 t + \Phi_2 t^2 + \Phi_3 t^3 + \Phi_4 t^4 + \Phi_5 t^5 + \dots$$

The first few formulas where terms of second order or greater in  $q_1$  and  $q_2$  are denoted by dots are as follows:

$$\begin{aligned} \Phi_0 &= -\mu - q_3 \\ \Phi_1 &= \frac{1}{\mu}(-2f_{20}q_1 + f_{11}q_2) + \dots \\ \Phi_2 &= \beta_2 + \frac{1}{\mu}(f_{20} + \alpha_2 f_{11}q_1 - 2\alpha_2 f_{02}q_2) + \frac{1}{\mu^2}(3f_{30}q_1 + f_{21}q_2) + \dots \\ \Phi_3 &= \beta_3 + \frac{1}{\mu}(\alpha_2 f_{11} - \alpha_3 f_{11}q_1 - 2\alpha_3 f_{02}q_2) + \frac{1}{\mu^2}(-f_{30} + 2\alpha_2 f_{21}q_1 + 2\alpha_2 f_{12}q_2) \\ &\quad + \frac{1}{\mu^3}(-4f_{40}q_1 - f_{31}q_2) + \dots \end{aligned}$$

Setting in these formulas  $q_1 = q_2 = 0$  we get the expressions of the Taylor coefficients of the organising centre  $g_k = \Phi_k|_{q_1=q_2=0}$  at a chord point  $m_0$ .

**2.1. Normal forms of the Minkowski set.** The following proposition together with explicit calculations from the normal forms prove theorem 1.4.

**Proposition 2.2.** *For a generic pair of  $M$  and  $N$  at any point  $q$  of a base chord  $(a_0, b_0)$  except the point  $b_0$  itself ( $\mu = 0$ ) the germ of the respective generating family  $\Phi$  is space-time contact equivalent to one of the standard versal deformations in parameters  $(\mu, q) \in \mathbb{R} \times \mathbb{R}^3$  of the function germs at the origin in the variable  $t$  of the type  $A_k$  for  $k = 1, \dots, 4$  as follows:*

$$\begin{aligned} A_1 : \Phi &= t^2 + \mu; & A_2 : \Phi &= t^3 + q_1 t + \mu; \\ A_3 : \Phi &= t^4 + q_2 t^2 + q_1 t + \mu; & A_4 : \Phi &= t^5 + q_3 t^3 + q_2 t^2 + q_1 t + \mu. \end{aligned}$$

**2.2. Recognition of transversal singularities.** If  $\beta_2$  is nonzero, that is the base tangent plane is not the osculating plane to the curve  $N$  then we always get a unique  $A_2$  singularity at the point  $\mu_c = -\frac{f_{20}}{\beta_2}$ . If however  $\beta_2 = 0$  then no caustic point occurs on the chord unless additionally  $f_{20} = 0$  in which case the whole chord is of type  $A_2$  and therefore belongs to the caustic. These are isolated chords and the situation when these occur at  $f_{20} = \beta_2 = 0$  is called the *flattening case*.

If the condition  $\beta_3 = \frac{f_{30}\beta_2^2}{f_{20}^2} + \frac{\alpha_2 f_{11}\beta_2}{f_{20}}$  holds then the caustic point at  $\mu_c$  will be of the type  $A_3$ . If in addition to the condition for an  $A_3$  singularity the condition  $\beta_4 = \frac{f_{40}\beta_2^3}{f_{20}^3} + \frac{\alpha_2 f_{21}\beta_2^2}{f_{20}^2} + \frac{\alpha_3 f_{11}\beta_2}{f_{20}} + \frac{\alpha_2^2 f_{02}\beta_2}{f_{20}}$  also holds, together with  $g_5$  being nonzero, then the caustic point at  $\mu_c$  will be of the type  $A_4$ . The singularities of this type occur at isolated points due to genericity.

In the flattening case  $f_{20} = \beta_2 = 0$ , in addition to the whole chord being of type  $A_2$ , there also exist two points where  $A_3$  singularities occur at  $\mu = \frac{-f_{11}\alpha_2 \pm \sqrt{f_{11}^2\alpha_2^2 + 4\beta_3 f_{30}}}{2\beta_3}$ .

### Proof of proposition 2.2

The proof of the proposition uses the property that  $\frac{\partial \Phi}{\partial \lambda} \neq 0$  which holds in the transversal case. The stability with respect to space-time contact equivalence of the germ  $\Phi$  with this property coincides with its stability with respect to standard contact equivalence. Therefore to show stability with respect to space-time contact equivalence we proceed by proving that each singularity in turn is generically versal with respect to standard contact equivalence, (see [2]).

Let  $\Phi : \mathbb{R} \times \mathbb{R} \times \mathbb{R}^3 \mapsto \mathbb{R}$  be an unfolding of a function  $g(t)$ ,  $t \in \mathbb{R}$  with parameters  $\mu, q \in \mathbb{R} \times \mathbb{R}^3$  and let  $g(t)$  have an  $A_k$  singularity at the origin.

Denote by  $\delta_{ij}$  the coefficients of the  $k$ -jet of  $\Phi$  at the origin where

$$\delta_{i1} = \frac{\partial^{i+1}\Phi}{\partial t^i \partial q_3}, \delta_{i2} = \frac{\partial^{i+1}\Phi}{\partial t^i \partial \mu}, \delta_{i3} = \frac{\partial^{i+1}\Phi}{\partial t^i \partial q_1} \text{ and } \delta_{i4} = \frac{\partial^{i+1}\Phi}{\partial t^i \partial q_2}.$$

The jet matrix for the family of functions  $\Phi$  shall be denoted  $M_4$  and let  $M_k$  with  $k \leq 4$  be the matrix consisting of the first  $k$  rows of  $M_4$ . We only consider  $k \leq 4$  due to genericity.

The matrix  $M_4 = (\delta_{ij})$  up to a factor of the rows for  $\mu$  nonzero is given by

$$M_4 = \begin{pmatrix} -1 & -1 & 0 & 0 \\ 0 & 0 & -2f_{20} & -f_{11} \\ 0 & \delta_{32} & \delta_{33} & \delta_{34} \\ 0 & \delta_{42} & \delta_{43} & \delta_{44} \end{pmatrix}$$

where

$$\begin{aligned} \delta_{32} &= -f_{20}, & \delta_{33} &= -\alpha_2 f_{11} \mu + 3f_{30}, & \delta_{34} &= -2\alpha_2 f_{02} \mu + f_{21}, \\ \delta_{42} &= -\alpha_2 f_{11} \mu + 2f_{30}, & \delta_{43} &= -\alpha_3 f_{11} \mu^2 + 2\alpha_2 f_{21} \mu - 4f_{40}, & \delta_{44} &= -2\alpha_3 f_{02} \mu^2 + 2\alpha_2 f_{12} \mu - f_{31}. \end{aligned}$$

Then function  $\Phi$  is right-versal if and only if the matrix  $M_k$  has rank  $k$ . Notice that the conditions  $g_1 = 0, \dots, g_k = 0$  define a Whitney stratification in the jet space of the embeddings. In fact each of these conditions outside  $\lambda$  being zero defines a regular hyper-surface in the space of germs and moreover those hyper-surfaces are mutually transversal since each equation  $g_i = 0$  involves only one variable  $\beta_i$  and can be solved for it in terms of coefficients  $f_{ji}$  and  $\alpha_s$ .

### Versality of an $A_1$ singularity

The proof is immediate because the matrix

$$M_1 = \begin{pmatrix} -1 & 1 & 0 & 0 \end{pmatrix}$$

always has the maximal rank of 1.  $\square$

### Versality of an $A_2$ singularity

The versality of the  $A_2$  singularities is determined by whether the matrix  $M_2$  has maximal rank 2 where

$$M_2 = \begin{pmatrix} -1 & 1 & 0 & 0 \\ 0 & 0 & -2f_{20} & -f_{11} \end{pmatrix}.$$

The matrix  $M_2$  has non-maximal rank only if both  $f_{11}$  and  $f_{20}$  vanish. If  $f_{20}$  vanishes and  $\beta_2 \neq 0$  then recall the caustic occurs at  $\mu = 0$  on the curve itself. If  $f_{20} = \beta_2 = 0$  then the whole chord belongs to the caustic. In this case the vanishing of  $\beta_2, f_{20}$  and  $f_{11}$  is non-generic. Therefore, away from the curve and surface,  $A_2$  singular points at  $\mu_c$  are versal.  $\square$

**Versality of an  $A_3$  singularity**

If  $f_{20} \neq 0$  then clearly the minor consisting of the first three columns of  $M_3$  as nonzero determinant. In the flattening case  $f_{20} = \beta_2 = 0$  the derivative matrix  $M_3$  takes the form

$$M_3 = \begin{pmatrix} -1 & 1 & 0 & 0 \\ 0 & 0 & 0 & f_{11} \\ 0 & 0 & -\mu\alpha_2 f_{11} + 3f_{30} & -2\mu\alpha_2 f_{02} + f_{21} \end{pmatrix}$$

and recall that the  $A_3$  singularities occur at  $\mu = \frac{-\alpha_2 f_{11} \pm \sqrt{\alpha_2^2 f_{11}^2 - 4\beta_3 f_{30}}}{2\beta_3}$ . For this value of  $\mu$  the determinant of both of the minors of  $M_3$  vanishes if either  $f_{11} = 0$  or  $\beta_3 = -\frac{4f_{11}^2 \alpha_2^2}{9f_{30}}$ . Neither of these conditions holds generically, so generically  $A_3$  singularities are versal.

**Versality of an  $A_4$  singularity.**

An  $A_4$  singularity occurs on the base chord at the point  $m_0$  when  $g_2 = g_3 = 0$  and

$$g_4 = \beta_4 + \frac{1}{\mu}(\alpha_3 f_{11} + \alpha_2^2 f_{02}) - \frac{1}{\mu^2} \alpha_2 f_{21} + \frac{1}{\mu^3} f_{40} = 0,$$

but  $g_5 \neq 0$ . Notice also that  $A_4$  cannot happen in the flattening case due to genericity.

The versality of  $A_4$  singularities holds if the determinant

$$\det(M_4) = 2f_{20}[\delta_{32}\delta_{44} - \delta_{34}\delta_{42}] - f_{11}[\delta_{32}\delta_{43} - \delta_{33}\delta_{42}]$$

is nonzero. Since generically  $A_4$  singularities cannot occur in the flattening case we assume that  $\beta_2$  and  $f_{20}$  are nonzero. The condition that the determinant is zero can be solved for  $f_{31}$  as a function of the other terms.

The codimension of the stratum which corresponds to an  $A_4$  singularity together with the vanishing of  $\det(M_4)$  is greater than 3 so  $A_4$  singularities are generically versal. This completes the proof of proposition 2.2. Explicit calculations from the normal forms completes the proof of theorem 1.4.  $\square$

**2.3. Propagating from the space curve in the tangential setting.** In the tangential case we use the same family as was used in the transversal case

$$\mathcal{F}_1(n, t, x, y, \mu, q) = \langle r_1(t) + \mu r_2(x, y) - q, n \rangle.$$

Here we assume that the base chord lies in the plane tangent to  $M$  at  $a_0$  (tangential setting). If the base chord and the tangent line to the curve  $N$  at  $b_0$  are not collinear then in some coordinate system  $(x, y, z)$  the base points  $a_0, b_0$  coincide with the points  $(0, 1, 0), (0, 0, 0)$ , the curve  $N$  at the origin is tangent to the  $x$ -axis and the tangent plane to the surface  $M$  coincides with the  $(x, y)$ -coordinate plane. Now the surface  $M$  is defined by the embedding

$$r_2 : U \rightarrow \mathbb{R}^3, \quad r_2(x, y) = (x, y + 1, f(x, y)), \quad (x, y) \in U \subset \mathbb{R}^2$$

where the function  $f(x, y)$  has zero 1-jet, and the curve  $N$  is defined by the embedding

$$r_1 : V \rightarrow \mathbb{R}^3, \quad r_1 : t \mapsto (t, \alpha(t), \beta(t)), \quad t \in V \subset \mathbb{R}$$

of some neighbourhood  $V$  of the origin in  $\mathbb{R}$  where  $\alpha(t)$  and  $\beta(t)$  start with second order terms.

After an appropriate stabilisation the initial generating family germ  $\mathcal{F}_1$  at the point  $\mu = \mu_0, t = 0, q = 0$  reduces to the form

$$(5) \quad \Phi(t, \varepsilon, q) = \beta(t) + (\mu_0 - \varepsilon) f \left( \frac{q_1 - t}{\mu_0 - \varepsilon}, \frac{\tilde{q}_2 - \alpha(t) - \varepsilon}{\mu_0 - \varepsilon} \right) - q_3$$

where  $\varepsilon = \mu_0 - \mu$  varies in the vicinity of the origin and  $\tilde{q}_2 = q_2 - \mu_0$ .



Consider the organising centre  $g(t, \varepsilon) = \Phi|_{q_1=\tilde{q}_2=q_3=0}$  of the family and decompose it as  $g(t, \varepsilon) = \sum_{i+j \geq 2} a_{ij} t^i \varepsilon^j$  where the first few terms are:

$$\begin{aligned} a_{20} &= \beta_2 + \frac{1}{\mu_0} f_{20}, & a_{11} &= -\frac{1}{\mu_0} f_{11}, & a_{02} &= \frac{1}{\mu_0} f_{02}, \\ a_{30} &= \beta_3 + \frac{1}{\mu_0} \alpha_2 f_{11} - \frac{1}{\mu_0^2} f_{30}, \\ a_{21} &= -\frac{2}{\mu_0} \alpha_2 f_{02} + \frac{1}{\mu_0^2} (f_{20} + f_{21}), \end{aligned}$$

The space-time contact equivalence of the families of type  $\Phi$  corresponds to *fibred contact equivalence* of the respective organising centres  $g(t, \varepsilon)$ : diffeomorphisms of the form  $(t, \varepsilon) \rightarrow (\hat{t}(t, \varepsilon), \hat{\varepsilon}(\varepsilon))$  and multiplications by nonzero functions act on  $g$ .

The well-known Arnold-Goryunov low dimensional fibred contact classification (which coincides with simple boundary classes) provides all generic space-time contact stable families depending on three parameters (here  $k = 2, 3$  or  $4$ ):

$$(6) \quad \begin{aligned} B_k &: \pm t^2 + \varepsilon^k + q_{k-2} \varepsilon^{k-2} + \dots + q_3, \\ C_2 &\approx B_2, \\ C_3 &: t^3 + t\varepsilon + q_1 \varepsilon + q_3, \\ C_3 &: t^4 + t\varepsilon + q_2 t^2 + q_1 \varepsilon + q_3, \\ F_4 &: t^3 + \varepsilon^2 + q_2 t \varepsilon + q_1 t + q_3. \end{aligned}$$

The proof of theorem 1.5 consists of checking the versality and genericity conditions for germs of the family  $\Phi$ .

Singularities of the type  $B_k$  occur  $a_{20} \neq 0$ . When the quadratic form of  $g(t, \varepsilon)$  is non-degenerate then the singularity of type  $B_2$  occurs. If the quadratic form is degenerate, that is  $4a_{20}a_{02} - a_{11}^2 = 0$ , the singularity is of type  $B_3$ . This occurs when  $\mu_0 = \frac{f_{11}^2 - 4f_{20}f_{02}}{4\beta_2 f_{02}}$  so every chord in the tangential setting has a singularity of type  $B_3$  (which may occur on the curve or at infinity). The  $B_3$  singularity can become more degenerate at isolated points to form the  $B_4$  type. This condition can be solved for  $\beta_3$  as a function of the other terms. Any further degenerations are excluded due to genericity.

The  $C_3$  singularity occurs when  $a_{20} = 0$  and both  $a_{20} \neq 0$  and  $a_{11} \neq 0$ . This happens at a single point on the chord when  $\mu_0 = -\frac{f_{20}}{\beta_2}$ . This can become more degenerate if  $a_{30} = 0$  and  $a_{40} \neq 0$  to form the singularity of type  $C_4$ . The singularity of type  $F_4$  belongs to the intersection of the closures of the  $B_3$  and  $C_3$  singularities and occurs when  $\mu_0 = -\frac{f_{20}}{\beta_2}$  and  $f_{11} = 0$ . Similar considerations using slightly different embeddings show that in the supertangential case, away from the curve and surface, only singularities of type  $B_2$  and  $C_3$  occur generically.

### 3. PROPAGATING FROM THE SURFACE IN THE TRANSVERSAL SETTING

We now turn our attention to the case where our initial starting manifold is a surface and the indicatrix of admissible velocities at each point defined by a space curve. Recall that in this case we use the generating function

$$\mathcal{F}_2(n, t, x, y, \lambda, q) = \langle \lambda r_1(t) + r_2(x, y) - q, n \rangle$$

where as before  $r_1(t)$  is the embedding with the image of the space curve  $N$ , and  $r_2(x, y)$  is the embedding with image the surface  $M$ . In this case we choose affine coordinates so that near a distinguished chord the surface is at the origin and the tangent plane is the  $(x, y)$ -coordinate

plane, and the space curve contains the point  $(0, 0, 1)$  and has tangent vector in the direction of the  $x$ -axis.

**Proposition 3.1.** *The family germ  $\mathcal{F}_2$  at a point corresponding to a point on the base chord is stably-equivalent to the product of the family germ  $\Phi(t, \mu, q) = \lambda(\beta(t) + 1) + [f(\hat{x}, \hat{y})] - q_3$  at the subset  $\widehat{S}_0 = \{t = x = y = 0, q_1 = q_2 = 0\}$  with a nonzero factor. Here we use the substitution  $\hat{x} = q_1 - \lambda t, \hat{y} = q_2 - \lambda \alpha(t)$ .*

**Proof.** Writing the family  $\mathcal{F}_2$  in the coordinate form we get

$$\mathcal{F} = An_1 + Bn_2 + Cn_3$$

where

$$\begin{aligned} A &= \lambda t + x - q_1 \\ B &= \lambda \alpha(t) + y - q_2 \\ C &= \lambda(\beta(t) + 1) + f(x, y) - q_3 \end{aligned}$$

As in the previous case we make an appropriate substitution, this time  $x = q_1 - \lambda t$  and  $y = q_2 - \lambda \alpha(t)$ , and use the Hadamard lemma to show that this is stably equivalent to the family

$$\Phi(t, \lambda, q) = \lambda\beta(t) + \lambda + [f(q_1 - \lambda t, q_2 - \lambda \alpha(t))] - q_3.$$

Expanding the function  $\Phi$  as a Taylor decomposition with respect to  $t$  at the origin where  $\Phi = \sum_{n=0}^{\infty} \Phi_k t^k$ , up to linear terms in  $q_1$  and  $q_2$ , has the first few coefficients:

$$\begin{aligned} \Phi_0 &= \lambda - q_3 \\ \Phi_1 &= -\lambda(2f_{20}q_1 + f_{11}q_2) \\ \Phi_2 &= \lambda(\beta_2 - f_{11}\alpha_2q_1 - 2f_{02}\alpha_2q_2 + \lambda f_{20} + 3\lambda f_{30}q_1 + f_{21}\lambda q_2) \end{aligned}$$

Setting in these formulas  $q_1 = q_2 = 0$  we get the expressions of the Taylor coefficients of the organising centre  $g_k = \Phi_k|_{q_1 = q_2 = 0}$  at a chord point  $m_0$ .

As with the propagating from the space curve case away from the initial starting manifold an  $A_2$  singularity occurs on the chord at  $\lambda_c = \frac{-\beta_2}{f_{20}}$ . This becomes more degenerate as an  $A_3$  singularity if additionally  $\beta_3 = \frac{\beta_2^2 f_{30}}{f_{20}^2} + \frac{\beta_2 f_{11} \alpha_2}{f_{20}}$  and type  $A_4$  if also  $\beta_4 = -\lambda^3 f_{40} + \lambda^2 \alpha_2 f_{21} - \lambda \alpha_3 f_{11} - \lambda \alpha_2^2 f_{02}$ . Notice that these conditions for the singularity to be more degenerate are the same as those in the propagating from the space curve case. In the flattening case when  $f_{20} = \beta_2 = 0$  the whole chord belongs to the caustic and is type  $A_2$  everywhere except two points  $\lambda = \frac{\alpha_2 f_{11} \pm \sqrt{\alpha_2^2 f_{11}^2 + 4\beta_3 f_{30}}}{2f_{30}}$  where singularities of type  $A_3$  occur.

Consider the derivative matrix given by

$$M_4 = \begin{pmatrix} -1 & 1 & 0 & 0 \\ 0 & 0 & -2f_{20} & -f_{11} \\ 0 & \delta_{32} & \delta_{33} & \delta_{34} \\ 0 & \delta_{42} & \delta_{43} & \delta_{44} \end{pmatrix}$$

where

$$\begin{aligned} \delta_{32} &= 2\lambda f_{20} + \beta_2, & \delta_{33} &= 3\lambda^2 f_{30} - \lambda \alpha_2 f_{11}, & \delta_{34} &= \lambda^2 f_{21} - 2\lambda \alpha_2 f_{02}, \\ \delta_{42} &= -3\lambda^2 f_{30} + 2\lambda \alpha_2 f_{11} + \beta_3, & \delta_{43} &= -4\lambda^3 f_{40} + 2\lambda^2 \alpha_2 f_{21} - \lambda \alpha_3 f_{11}, \\ \delta_{44} &= -\lambda^3 f_{31} + 2\lambda^2 \alpha_2 f_{12} - 2\lambda \alpha_3 f_{02}. \end{aligned}$$

We can use the same arguments as we used in the propagating from the space curve case to show that all the generic  $A_k$  singularities are versal.

**3.1. Propagating from the surface in the tangential setting.** Assume that the base chord is in the tangential setting, that is it lies in the plane tangent to  $M$  at  $a_0$  but it is not collinear with the tangent line to the curve  $N$  at  $b_0$ . In some coordinate system the surface contains the origin at  $a_0$  and the tangent plane at this point is the  $(x, y)$ -coordinate plane, the curve passes through the point  $(0, 1, 0)$  at  $b_0$  and the tangent line has the same direction as the  $x$ -axis. Using appropriate embeddings  $r_1(t)$  and  $r_2(x, y)$  the generating  $\mathcal{F}_2$  can be expanded as a vector to give

$$\mathcal{F}_2 = (x + \lambda t - q_1)n_1 + (y + \lambda(\alpha(t) + 1) - q_2)n_2 + ((f(x, y) + \lambda\beta(t) - q_3)n_3.$$

**Proposition 3.2.** *Using the substitution  $x = q_1 - \lambda t$ ,  $y = q_2 - \lambda(\alpha(t) + 1)$  the family  $\mathcal{F}_2$  at the point  $\lambda = \lambda_0, t = 0, q = 0$  is stably equivalent to the family*

$$(7) \quad \Phi(t, \varepsilon, q) = (\lambda_0 + \varepsilon)\beta(t) + f(q_1 - (\lambda_0 + \varepsilon)t, \tilde{q}_2 - \lambda_0\alpha(t) - \varepsilon\alpha(t) - \varepsilon).$$

where  $\lambda = (\lambda_0 + \varepsilon)$  varies in the vicinity of the origin and  $q_2 = \tilde{q}_2 + \lambda_0$ .

Consider the organising centre  $g(t, \varepsilon) = \Phi|_{q_1=\tilde{q}_2=q_3=0}$  of the family and decompose it as  $g(t, \varepsilon) = \sum_{i+j \geq 2} a_{ij}t^i\varepsilon^j$  where the first few terms are

$$\begin{aligned} a_{20} &= f_{20}\lambda_0^2 + \lambda_0\beta_2, & a_{11} &= f_{11}\lambda_0, & a_{02} &= f_{02}, \\ a_{30} &= \lambda_0\beta_3 + f_{11}\lambda_0^2\alpha_2 - f_{30}\lambda_0^3, \\ a_{21} &= \beta_2 + 2f_{02}\lambda_0\alpha_2 - f_{21}\lambda_0^2 + 2f_{20}\lambda_0, \\ a_{12} &= f_{11} - f_{12}\lambda_0, & a_{03} &= -f_{03}. \end{aligned}$$

The list of generic singularities coincides with the list (6).

The whole of each chord in the tangential setting is of type at  $B_2$  except for at most two points which can be more degenerate. Generically each chord will consist of a  $B_3$  singularity and a  $C_3$  singularity. At isolated chords one of these can be more degenerate to form either  $B_4$  or  $C_4$ . Also at isolated chords the  $B_3$  and  $C_3$  singularities can occur at the same point to give an  $F_4$  singularity.

When the quadratic form is degenerate, that is  $4a_{20}a_{02} - a_{11}^2 = 0$ , a singularity of type  $B_3$  occurs. This happens at  $\lambda = \frac{4f_{02}\beta_2}{f_{11}^2 - 4f_{02}f_{20}}$ . For isolated chords one of these can be more degenerate giving the singularity of type  $B_4$ .

This condition can be solved for  $f_{03}$  as a function of the other terms.

Singularities of type  $C_3$  occur at  $\lambda = -\frac{\beta_2}{f_{20}}$ . This can be more degenerate to form a  $C_4$  singularity if  $\beta_3 = \frac{\beta_2^2 f_{30}}{f_{20}^2} + \frac{\beta_2 f_{11} \alpha_2}{f_{20}}$ . An  $F_4$  singularity will result if both the conditions for a  $C_3$  and a  $B_3$  singularity occur, namely if  $\lambda = -\frac{\beta_2}{f_{20}}$  and  $f_{11} = 0$ . Further degenerations are excluded due to genericity.

Checking the versality and genericity conditions for germs of the family  $\Phi$  completes the proof of Theorem 1.8.

Similar considerations using different embeddings show that in the supertangential case, away from the curve and surface, generically only singularities of type  $B_2$  and  $C_3$  occur.  $\square$

**3.2. Propagating from the surface in the transversal case in the vicinity of the surface.** Up until now we have assumed that  $\lambda$  is nonzero and have classified the singularities away from the surface and space curve. In this section we study the generic caustic near the surface itself, that is when  $\lambda$  is close to zero. We use the standard generating family in the propagating from the surface case  $\mathcal{F}_2$  and proposition 3.1 implies the generating family is stably equivalent to

$$\Phi(t, \lambda, q) = \lambda(\beta(t) + 1) + f(q_1 - \lambda t, q_2 - \lambda\alpha(t)) - q_3.$$

This can be written as

$$\Phi = \lambda \left( \frac{f - f_0}{\lambda} + \beta + 1 \right) + f_0 - q_3$$

where  $f_0 = f|_{q_1=q_2=0}$  and  $\frac{f-f_0}{\lambda}$  is smooth. Introduce the new parameter  $\tilde{q}_3 = -q_3 + f_0$  which vanishes on the surface, yielding

$$\Phi = \lambda \left( \frac{f - f_0}{\lambda} + \beta + 1 \right) + \tilde{q}_3.$$

Denote by  $\Phi_0, \Phi_1, \dots$  the terms of the power series decomposition in  $\lambda$  of the contents of the brackets. With terms of order greater than 4 in  $t$  or greater than 1 in  $q_1$  and  $q_2$  denoted by dots the generating function  $\Phi$  is written

$$(8) \quad \mathcal{F} = \tilde{q}_3 + \lambda (\Phi_0 + \dots + \lambda(\Phi_1 + \dots) + \lambda^2(\Phi_2 + \dots) + \dots)$$

where

$$\Phi_0 = 1 + \beta(t) + (-2f_{02}\alpha(t) - f_{11}t)q_2 + (-2f_{20}t - f_{11}\alpha(t))q_1$$

and

$$\begin{aligned} \Phi_1 = & f_{11}t\alpha(t) + f_{20}t^2 + f_{02}\alpha(t)^2 + \\ & \left( 2f_{12}t\alpha(t) + f_{21}t^2 + 3f_{03}\alpha(t)^2 \right) q_2 + \left( 2f_{21}t\alpha(t) + f_{12}\alpha(t)^2 + 3f_{30}t^2 \right) q_1. \end{aligned}$$

**Proposition 3.3.** *The family germ  $\mathcal{F}$  can be written in the form  $\mathcal{F} = \lambda + \tilde{q}_3 H(t, q_1, q_2, \tilde{q}_3)$  where the lower degree terms with respect to  $\tilde{q}_3$  of function  $H$  are:  $H = \frac{1}{\Phi_0} + \frac{\Phi_1 \tilde{q}_3}{\Phi_0^2} + \dots$*

**Lemma 3.4.** *Assume  $H(t, q_1, q_2, \tilde{q}_3)$  is  $\mathcal{R}^+$ -versal with respect to  $q_1$  and  $q_2$  only; then the family germ  $\mathcal{F} = \lambda + \tilde{q}_3 H$  is space-time stable with respect to deformations inside the space  $W = \lambda + \tilde{q}_3 \tilde{H}(t, q_1, q_2, \tilde{q}_3)$  such that  $\frac{\partial W}{\partial \tilde{q}_3} \neq 0$ .*

**Proposition 3.5.** *For generic curve and surface germs in the transversal setting the function  $H(t, \lambda, q)$  is versal for standard  $\mathcal{R}^+$ -equivalence with respect to  $q_1$  and  $q_2$  only.*

The first few terms of the Taylor decomposition of  $H$  with respect to  $t$  at the origin, namely  $H = \sum_{k=0} H_k(t, q)t^k$ , up to first order terms in  $q_i$ , are as follows.

$$\begin{aligned} H_0 &= 1, \\ H_1 &= 2f_{20}q_1 + f_{11}q_2, \\ H_2 &= -\beta_2 + f_{11}\alpha_2q_1 + 2f_{02}\alpha_2q_2 + f_{20}\tilde{q}_3, \\ H_3 &= -\beta_3 + (-4\beta_2f_{20} + f_{11}\alpha_3)q_1 + (-2\beta_2f_{11} + 2f_{02}\alpha_3)q_2 + f_{11}\alpha_2\tilde{q}_3, \end{aligned}$$

Setting in these formulas  $q_1 = q_2 = \tilde{q}_3 = 0$  we get the following expressions of the Taylor coefficients of the organising centre  $h_k = H_k|_{q_1=q_2=\tilde{q}_3=0}$ :

$$h_0 = 1, \quad h_1 = 0, \quad h_2 = -\beta_2, \quad h_3 = -\beta_3, \quad h_4 = \beta_2^2 - \beta_4$$

The function  $H$  has a singularity of type  $A_2$  if  $\beta_2 = 0$  and  $\beta_3 \neq 0$ . If  $\beta_2 = \beta_3 = 0$  and  $\beta_4 \neq 0$  then the function  $H$  has a singularity of type  $A_3$ . More degenerate singularities are excluded due to genericity.

In order for  $H$  to be  $\mathcal{R}^+$ -versal with respect to  $q_1$  and  $q_2$  only at an  $A_k$  singularity for  $k = 2, 3$  we need the first  $k - 1$  rows of the jet matrix

$$M_{k-1} = \begin{pmatrix} \frac{\partial^2 H}{\partial q_1 \partial t} & \frac{\partial^2 H}{\partial q_2 \partial t} \\ \frac{\partial^3 H}{\partial q_1 \partial t^2} & \frac{\partial^3 H}{\partial q_2 \partial t^2} \end{pmatrix}$$

to have maximal rank  $k - 1$ .

#### Versality of $A_2$ singularities on the surface $M$

An  $A_2$  singularity occurs if  $\beta_2 = 0$  and  $\beta_3 \neq 0$ . Recall that this is the necessary and sufficient condition that the tangent plane to the curve is the osculating plane with 3 point contact. The  $A_2$  singularities are versal if the matrix

$$M_1 = \begin{pmatrix} 2f_{20} & f_{11} \end{pmatrix}$$

has rank 1.

Clearly the vanishing of  $\beta_2, f_{20}$  and  $f_{11}$  provide a set of non-generic conditions so  $A_2$  singularities in the vicinity of the surface are versal.

#### Versality of $A_3$ singularities on the surface $M$

An  $A_3$  singularity occurs if  $\beta_2 = 0, \beta_3 = 0$  and  $\beta_4 \neq 0$ . This is the condition that the tangent plane to curve is the osculating plane and has 4 point contact (at a torsion zero). The  $A_3$  singularities are versal if the matrix

$$M_2 = \begin{pmatrix} 2f_{20} & f_{11} \\ \alpha_2 f_{11} & 2\alpha_2 f_{02} \end{pmatrix}$$

has rank 2. The condition  $\det(M_2) = 0$  together with the necessary conditions  $\beta_2 = \beta_3 = 0$  singularity provide a non-generic condition so  $A_3$  singularities are versal.

Since the generic singularities of the function  $H$  are  $\mathcal{R}^+$ -versal, lemma 3.4 implies that the generic singularities of the function  $\mathcal{F}$  are space-time stable inside the space  $W$ .

At  $A_2$  type points on the surface the caustic is smooth and transversally intersects the surface  $M$ . The respective generating family germ is space-time equivalent to the normal form:

$$\mathcal{F} = \lambda + \tilde{q}_3(t^3 + q_1 t + 1)$$

At  $A_3$  type points on the surface the caustic has a cuspidal edge that transversally intersects the surface  $M$ . The respective generating family germ is space-time equivalent to the normal form:

$$\mathcal{F} = \lambda + \tilde{q}_3(t^4 + q_1 t^2 + q_2 t + 1). \quad \square$$

**3.3. Propagating from the surface in the Tangential Case in the vicinity of the surface.** In this case the caustic is space-time contact equivalent to one of the following normal forms (see [6]):

$$\begin{aligned} \widehat{B}_2 : \quad & \tilde{q}_3 + \lambda(t^2 + q_1); & \widehat{B}_3 : \quad & \tilde{q}_3 + \lambda(t^2 \pm \lambda^2 + \lambda q_1 + q_2); \\ \widehat{C}_3 : \quad & \tilde{q}_3 + \lambda(t^3 + \lambda t + \lambda + q_1 t + q_2). \end{aligned}$$

The caustic at a  $\widehat{B}_2$  singularity consists only of the surface  $M$  and the criminant is a smooth surface with first order tangency with the surface  $M$ . At a  $\widehat{B}_3$  singularity the criminant is diffeomorphic to a semi-cubic cylinder and has second order tangency with the surface  $M$  at  $a_0$  (see figure 1). At a  $\widehat{C}_3$  singularity the criminant is diffeomorphic to a folded Whitney umbrella and the caustic is a smooth surface. The cuspidal edge of the folded Whitney umbrella has first

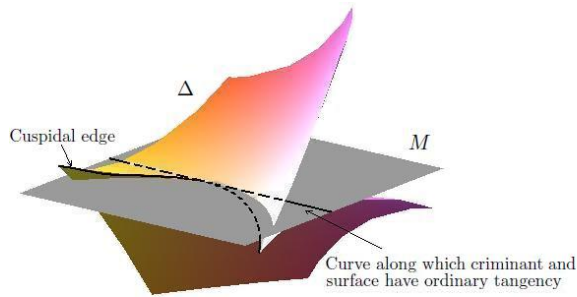


FIGURE 1. The envelope  $B(M, N)$  at a  $\widehat{B}_3$  singularity near the surface  $M$  (plane in figure). Here the caustic is empty and the crumpled  $\Delta$  is a cuspidal edge with second order tangency with the surface at  $a_0$ .

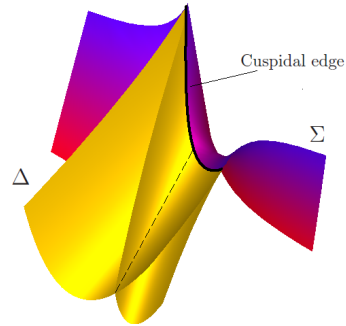


FIGURE 2. The caustic and the crumpled shown together at a  $\widehat{C}_3$  singularity. Here the crumpled  $\Delta$  is a folded Whitney umbrella and the caustic  $\Sigma$  is a smooth surface. The cuspidal edge of  $\Delta$  has third order tangency with  $\Sigma$ .

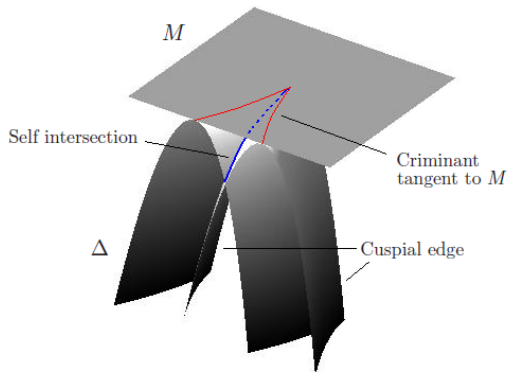


FIGURE 3. The crumpled and the surface  $M$  (plane in figure) are shown together at a  $\widehat{C}_3$  singularity. Here the crumpled and  $M$  have ordinary tangency along a cusp.

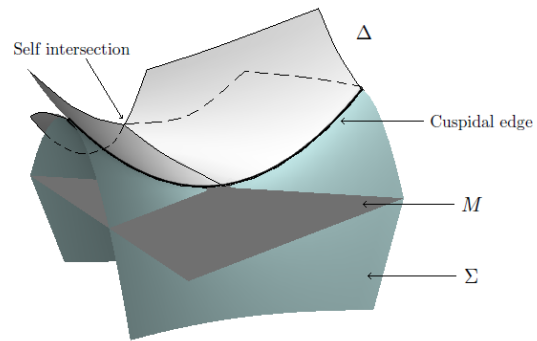


FIGURE 4. The envelope  $B(M, N)$  at a  $\widehat{C}_3$  singularity near the surface  $M$  (plane in figure). Here the crumpled  $\Delta$  is a folded Whitney umbrella and the caustic  $\Sigma$  is a smooth surface.

order tangency with the surface  $M$  at  $a_0$  and third order tangency with the caustic (see figure 2). Two additional views are shown in figures 3 and 4.

#### REFERENCES

- [1] A.A. Agrachev, *Methods of control theory in non-holonomic geometry*, Proc. Int. Congr. Math., Zürich, 1994, Birkhäuser, Basel, (1995), 12 -19.
- [2] J.W. Bruce and P.J. Giblin, *Curves and Singularities* second edition, Cambridge University Press (1992).
- [3] P.J. Giblin and V.M. Zakalyukin, *Singularities of centre symmetry sets*, London Math Soc. **90** (3), (2005), 132-166.
- [4] P.J. Giblin and V.M. Zakalyukin, *Recognition of centre symmetry set singularities*, Geom. Dedicata **130**, (2007), 43-58. DOI: [10.1007/s10711-007-9204-2](https://doi.org/10.1007/s10711-007-9204-2)

- [5] P.J. Giblin, J.P. Warder and V.M. Zakalyukin, *Bifurcations of affine equidistants*, Proc. Steklov Inst. Math., **269**, (2009), 78-96.
- [6] G.M. Reeve, *Singularities of systems of chords in affine space*, PhD thesis, University of Liverpool, (2012).
- [7] G.M. Reeve and V.M. Zakalyukin, *Singularities of the Minkowski set and affine equidistants for a Curve and a Surface*, Topology and its Applications, **159** (2), (2012), 555-561. DOI: [10.1016/j.topol.2011.09.031](https://doi.org/10.1016/j.topol.2011.09.031)
- [8] L.P. Stunzhas, *Local singularities of chord sets*, Mat. Zametki **83** (2), (2008), 286-304.
- [9] V.M. Zakalyukin *Singularities of caustics in generic translation-invariant control problems*, Journal of Mathematical Sciences, **126** (4), (2005), 1354-1360